

# Finite-volume effects and the electromagnetic contributions to kaon and pion masses

Claude Bernard  
Washington University  
Saint Louis, Missouri, USA

Co-authors: S. Basak, A. Bazavov, C. DeTar, E. Freeland, J. Foley, Steven  
Gottlieb, U.M. Heller, J. Laiho, L. Levkova, M. Oktay, J. Osborn , R.L. Sugar,  
A. Torok, D. Toussaint, R.S. Van de Water, R. Zhou

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# Motivation

- ◆ Disentangling electromagnetic and isospin-violating effects in the pions and kaons is long-standing issue.
- ◆ Crucial for determining light-quark masses.
  - Fundamental parameters in Standard Model; important for phenomenology.
  - Size of EM contributions is largest uncertainty in determination of  $m_u/m_d$ .

	$m_u$ [GeV]	$m_d$ [GeV]	$m_u/m_d$
value	1.9	4.6	0.42
statistics	0.0	0.0	0.00
lattice	0.1	0.2	0.01
perturbative	0.1	0.2	--
<b>EM</b>	<b>0.1</b>	<b>0.1</b>	<b>0.04</b>

MILC,  
[arXiv:0903.3598](https://arxiv.org/abs/0903.3598)

- Reduce error by calculating EM effects on the lattice.

# Background

- ◆ EM error in  $m_u/m_d$  dominated by error in  $(M_{K^+}^2 - M_{K^0}^2)^\gamma$ , where  $\gamma$  indicates the EM contribution.
- ◆ Dashen (1960) showed that at leading order EM splittings are mass independent:

$$(M_{K^+}^2 - M_{K^0}^2)^\gamma = (M_{\pi^+}^2 - M_{\pi^0}^2)^\gamma$$

- ◆ Parameterize higher order effects (“corrections to Dashen’s theorem”) by

$$(M_{K^+}^2 - M_{K^0}^2)^\gamma = (1 + \epsilon)(M_{\pi^+}^2 - M_{\pi^0}^2)^\gamma$$

- Note:  $\epsilon$  is not exactly same as quantity defined by FLAG ([Colangelo et al., arXiv:1310.8555](#)), which uses experimental pion splittings. But EM splitting should be  $\approx$  experimental splitting, since isospin violations for pions are small. Using the experimental splitting gives an alternative result, which enters systematic error estimate.

# Ensembles

- ◆ Table of ensembles used in the analysis:

$\approx a[\text{fm}]$	Volume	$\beta$	$m_l/m_s$	# configs.	$L$ (fm)	$m_\pi L$
0.12	$12^3 \times 64$	6.76	0.01/0.05	1000	1.4	2.7
	$16^3 \times 64$	6.76	0.01/0.05	1003	1.8	3.6
	$20^3 \times 64$	6.76	0.01/0.05	2254	2.3	4.5
	$28^3 \times 64$	6.76	0.01/0.05	274	3.2	6.3
	$20^3 \times 64$	6.76	0.007/0.05	1261	2.3	3.8
	$24^3 \times 64$	6.76	0.005/0.05	2099	2.7	3.8
0.09	$28^3 \times 96$	7.09	0.0062/0.031	1930	2.3	4.1
	$40^3 \times 96$	7.08	0.0031/0.031	1015	3.3	4.2
0.06	$48^3 \times 144$	7.47	0.0036/0.018	670	2.8	4.5

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- From [Bijnens and Daniellson \[PRD 75, 104505 \('07\)\]](#), quenched QED is sufficient for a controlled calculation of  $\epsilon$  at NLO in SU(3) ChPT.

# Ensembles

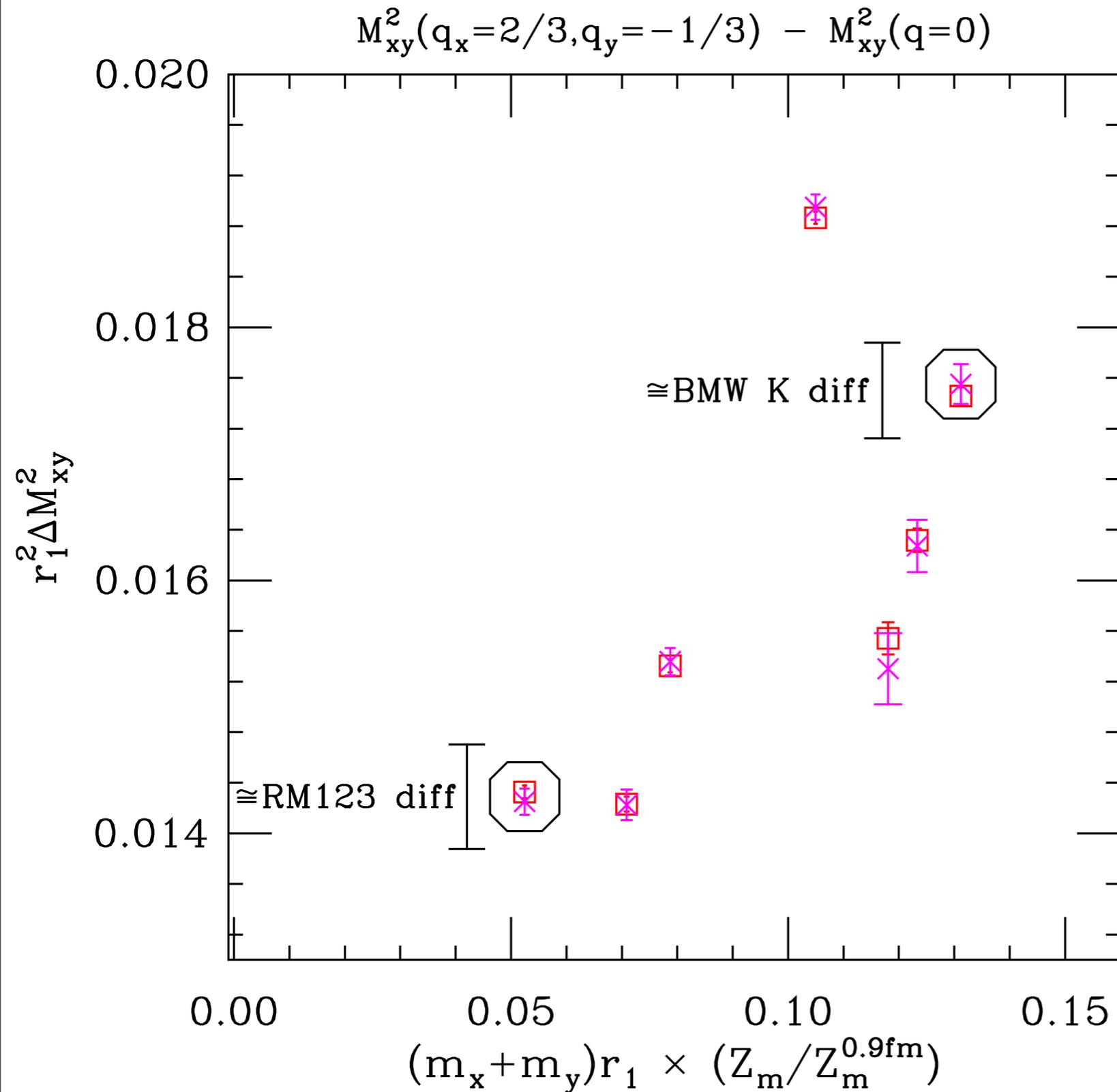
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◆ These are dynamical QCD ( $N_F=3$ , asqtad) ensembles, with quenched, noncompact QED.

- From [Bijnens and Daniellson \[PRD 75, 104505 \('07\)\]](#), quenched QED is sufficient for a controlled calculation of  $\epsilon$  at NLO in SU(3) ChPT.
- Small volumes used only to test our understanding of finite-volume effects, not for final analysis.

# Finite-Volume Effects



- Difference between  $20^3$  ( $\square$ ) and  $28^3$  ( $\times$ ) ensembles at  $a \approx 0.12$  fm is small compared to what we expect from BMW [arXiv:1201.2787], and RM123 [arXiv:1303.4896] results.
- We are not currently able to resolve the differences (consistent with zero).
  - Sign of the difference actually varies fairly randomly as quark masses change.
- Our recent work has been focused on understanding the (surprisingly small) FV effects in our data.

# Finite-Volume Effects in ChPT

- ◆ Hayakawa and Uno [arXiv:0804.2044] calculated the EM finite-volume effects in ChPT.
  - Use noncompact realization of QED on the lattice, as we do.
  - Found rather large effects.
  - But noncompact QED in finite-volume is not uniquely defined:
    - It is necessary to drop some zero modes, but dropping others appears to be optional.
    - In Coulomb gauge, action for  $A_0$  is:  $\frac{1}{2} \int (\partial_i A_0)^2$ .
      - For path integral to be convergent, need to drop  $A_0$  modes for 3-momentum  $\vec{k}=0$ , any  $k_0$ .
    - Action for  $A_i$  is:  $\frac{1}{2} \int \left[ (\partial_0 A_i)^2 + (\partial_j A_i)^2 \right]$ .
      - Here, only required to drop mode with 4-momentum  $k_\mu=0$ .
      - Hayakawa & Uno drop all  $A_i$  modes with  $\vec{k}=0$ .
      - MILC keeps modes with  $\vec{k}=0$ ,  $k_0 \neq 0$ .

# Finite-Volume Coulomb-Gauge Propagator

$$\langle A_0(k) A_0(-k) \rangle = \begin{cases} \frac{1}{\vec{k}^2}, & \vec{k} \neq 0; \\ 0, & \vec{k} = 0. \end{cases}$$
$$\langle A_i(k) A_j(-k) \rangle = \begin{cases} \frac{1}{k^2} \left( \delta_{ij} - \frac{k_i k_j}{\vec{k}^2} \right), & \vec{k} \neq 0; \\ 0, & \vec{k} = 0. \end{cases}$$

Hayakawa-Uno

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MILC

# Finite-Volume Coulomb-Gauge Propagator

- ◆ Hayakawa and Uno have an argument for dropping zero modes based on the problem of having a single electric charge on a torus, due to Gauss's law.
  - Gauss's law comes from the equation of motion for  $A_0$ .
  - Hayakawa & Uno and MILC drop the same modes for  $A_0$  so Gauss's law solution is the same for both.
  - Difference is only for  $\vec{k}=0$  modes for  $A_i$ .

# Chiral Perturbation Theory

◆ Staggered version of NLO SU(3)  $\chi$ PT [C.B. & Freeland, arXiv:1011.3994]:

$$\begin{aligned}\Delta M_{xy,5}^2 &= q_{xy}^2 \delta_{EM} - \frac{1}{16\pi^2} e^2 q_{xy}^2 M_{xy,5}^2 [3 \ln(M_{xy,5}^2 / \Lambda_\chi^2) - 4] \\ &\quad - \frac{2\delta_{EM}}{16\pi^2 f^2} \frac{1}{16} \sum_{\sigma,\xi} [q_{x\sigma} q_{xy} M_{x\sigma,\xi}^2 \ln(M_{x\sigma,\xi}^2) - q_{y\sigma} q_{xy} M_{y\sigma,\xi}^2 \ln(M_{y\sigma,\xi}^2)] \\ &\quad + c_1 q_{xy}^2 a^2 + c_2 q_{xy}^2 (2m_\ell + m_s) + c_3 (q_x^2 + q_y^2) (m_x + m_y) + c_4 q_{xy}^2 (m_x + m_y) + c_5 (q_x^2 m_x + q_y^2 m_y)\end{aligned}$$

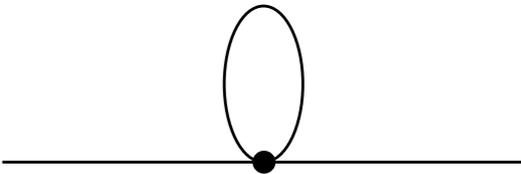
- x,y are the valence quarks.
- $q_x, q_y$  are **quark** charges;  $q_{xy} \equiv q_x - q_y$  is **meson** charge.
- $\delta_{EM}$  is the LO LEC;  $\xi$  is the staggered taste
- $\sigma$  runs over sea quarks ( $m_u, m_d, m_s$ , with  $m_u = m_d \equiv m_\ell$ )
- Finite-volume corrections coming from the sunset and photon tadpole graphs are non-trivial.
  - (FV corrections to meson tadpole are known from standard ChPT and are quite small).

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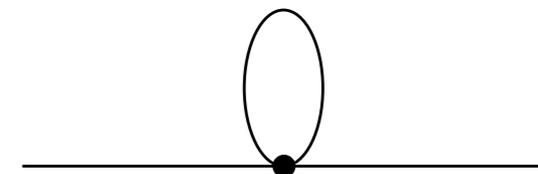
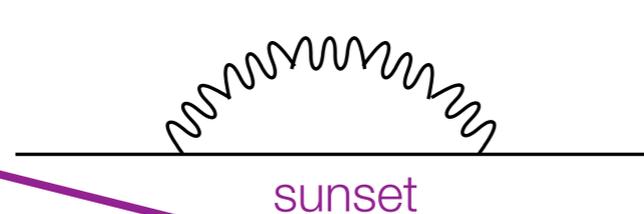
meson EM tadpole  
(from short-distance photons)

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long-distance  
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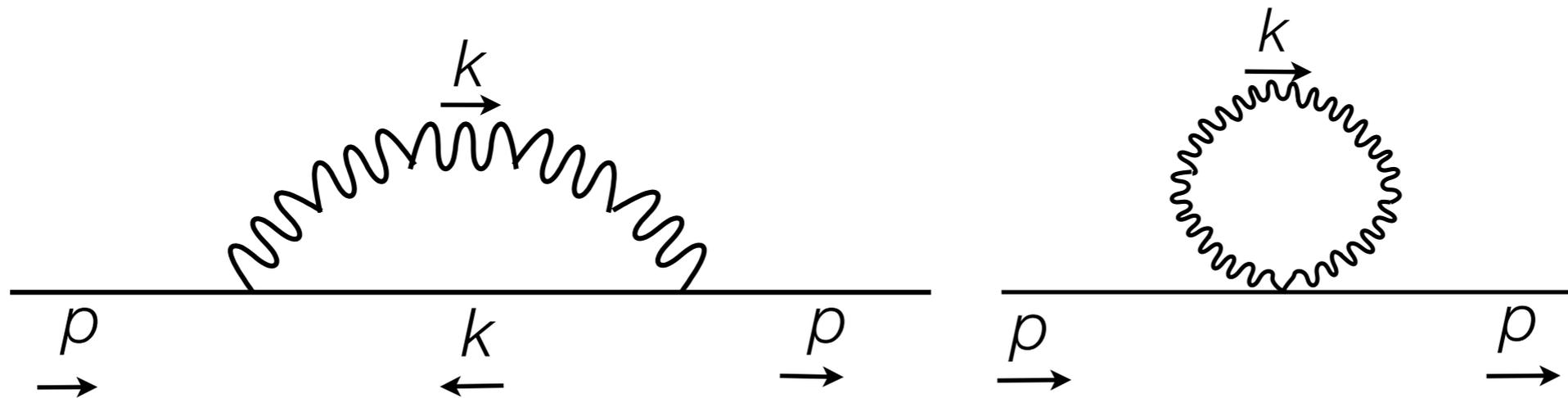
$$- \frac{2\delta_{EM}}{16\pi^2 f^2} \frac{1}{16} \sum_{\sigma,\xi} [q_{x\sigma} q_{xy} M_{x\sigma,\xi}^2 \ln(M_{x\sigma,\xi}^2) - q_{y\sigma} q_{xy} M_{y\sigma,\xi}^2 \ln(M_{y\sigma,\xi}^2)]$$

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# Finite-Volume ChPT

- ◆ Need to add photon diagrams together in order for Coulomb-gauge finite-volume difference (FV -  $\infty V$ ) to be well-defined.

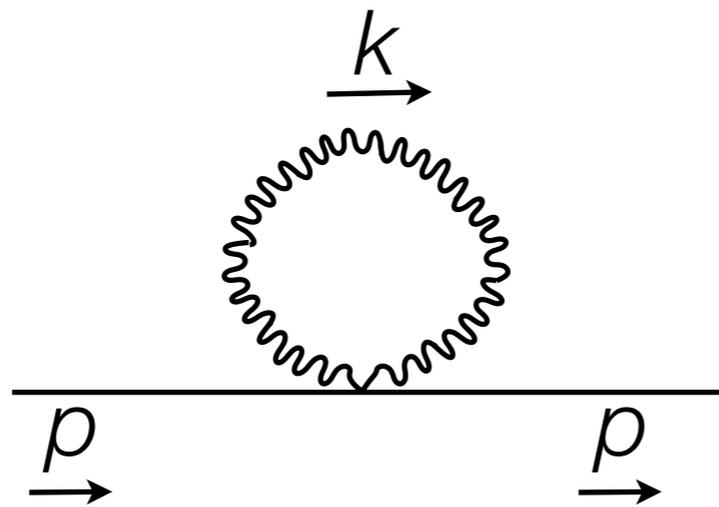


- ◆ Can then perform brute force difference of FV sum (over  $2\pi n_i/L$  and  $2\pi n_0/T$ ) from  $\infty V$  integral.

# Evaluation of FV difference

- ◆ Evaluate difference of sum and integral by VEGAS.
- ◆ Take VEGAS integrand as difference between  $\infty V$  integrand, and its evaluation at weighted average of the 16 corners of the FV hypercube containing the point.
- ◆ Checked against [Hayakawa-Uno](#) result (written in terms of 1-d integral over special functions).

# Photon Tadpole Graph



◆ There is a difference in FV part of photon tadpole between **Hayakawa-Uno (HU)** and **MILC** when  $\vec{k} = 0$ :

- **HU** omits the  $\vec{k} = 0$  piece entirely.

- For **MILC**, FV integrand is  $\frac{3}{k^2} = \frac{3}{k_0^2}$ , as long as  $k_0 \neq 0$ .

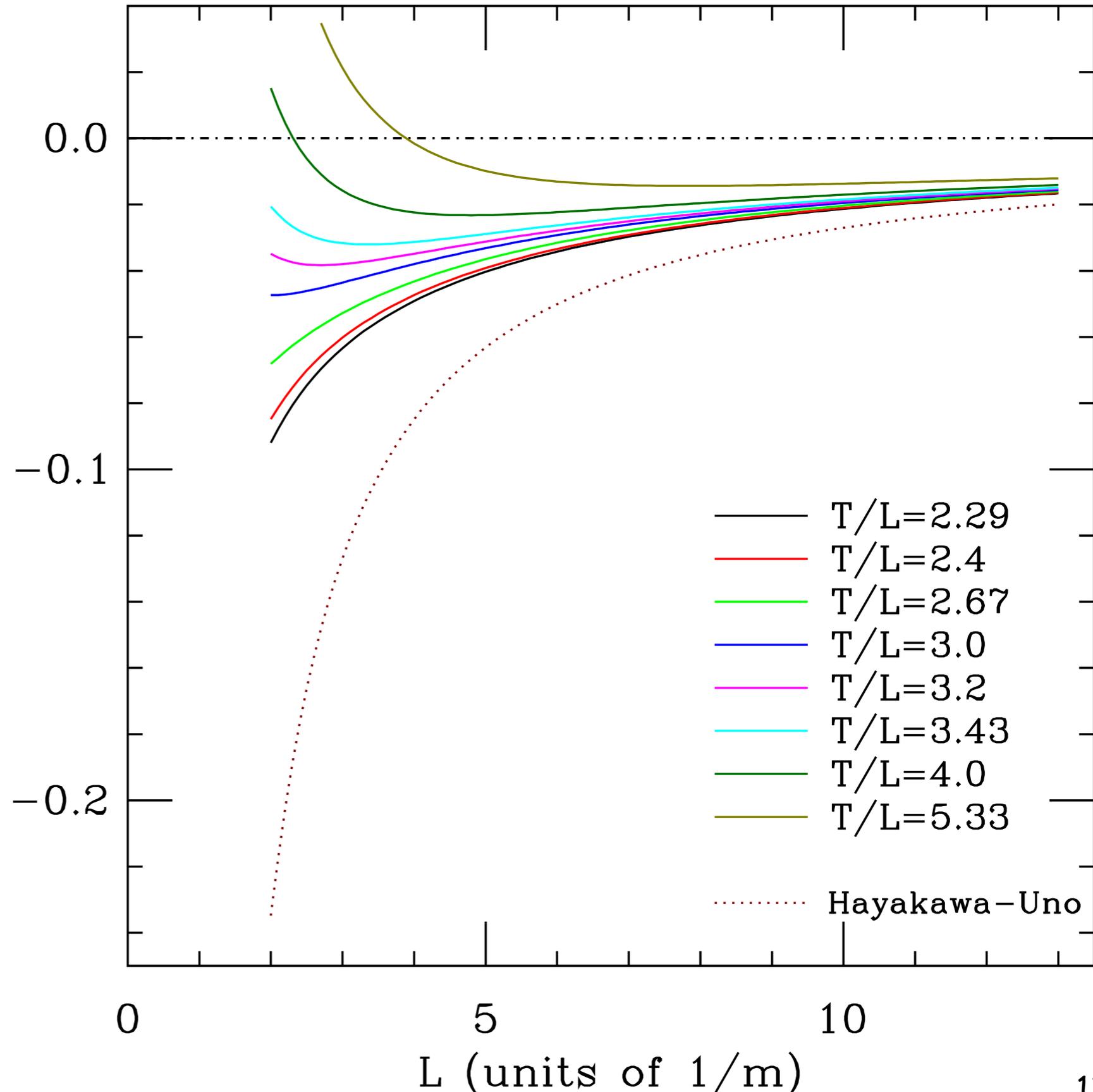
- Difference (**MILC-HU**) =  $\frac{q^2}{L^3 T} \sum_{n_0 \neq 0} \frac{3}{(2\pi n_0/T)^2} = \frac{q^2 T}{4L^3}$ .

- Our formulation has subtle  $T, L$  dependence.

- Fine if  $L \rightarrow \infty$  first, or if both  $T, L \rightarrow \infty$  with fixed ratio, but not if  $T \rightarrow \infty$  first.

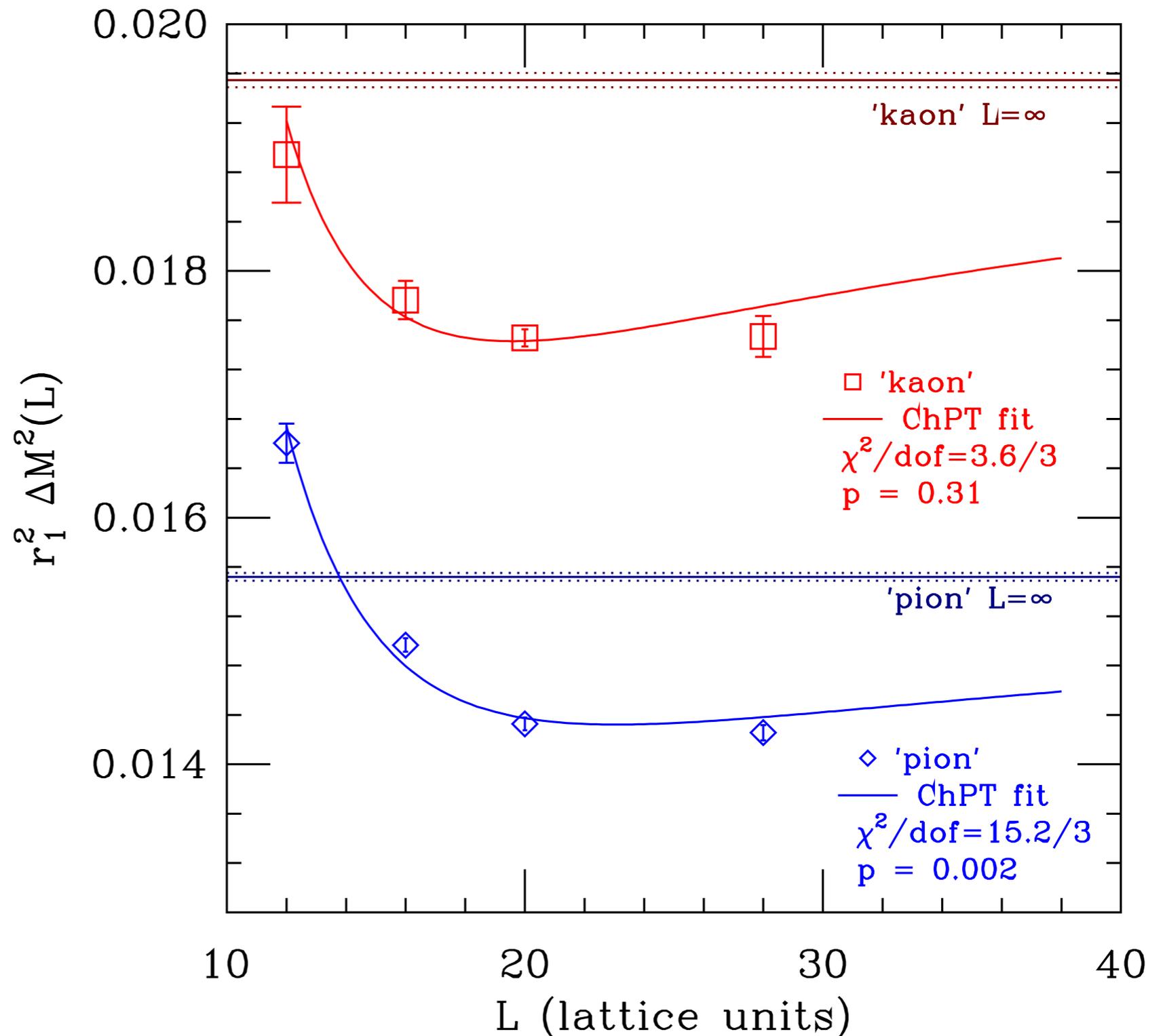
# Finite-Volume Corrections

- Comparison of MILC and H-U FV corrections.
  - ♦ An overall factor of  $e^2 m^2$ , (where  $e$  and  $m$  are charge & mass of the meson) has been taken out.
- $T/L$  values are the ones of our lattices.
  - ♦  $T/L = 4.0, 5.33$  are the small lattices ( $\sim 1.4$  fm,  $\sim 1.8$  fm) used only for investigating FV effects.
- H-U results are insensitive to  $T$  in this range. (In their paper, they calculate in the  $T = \infty$  limit only.)
- Our FV corrections are a factor of 2-3 less in most of the relevant range!



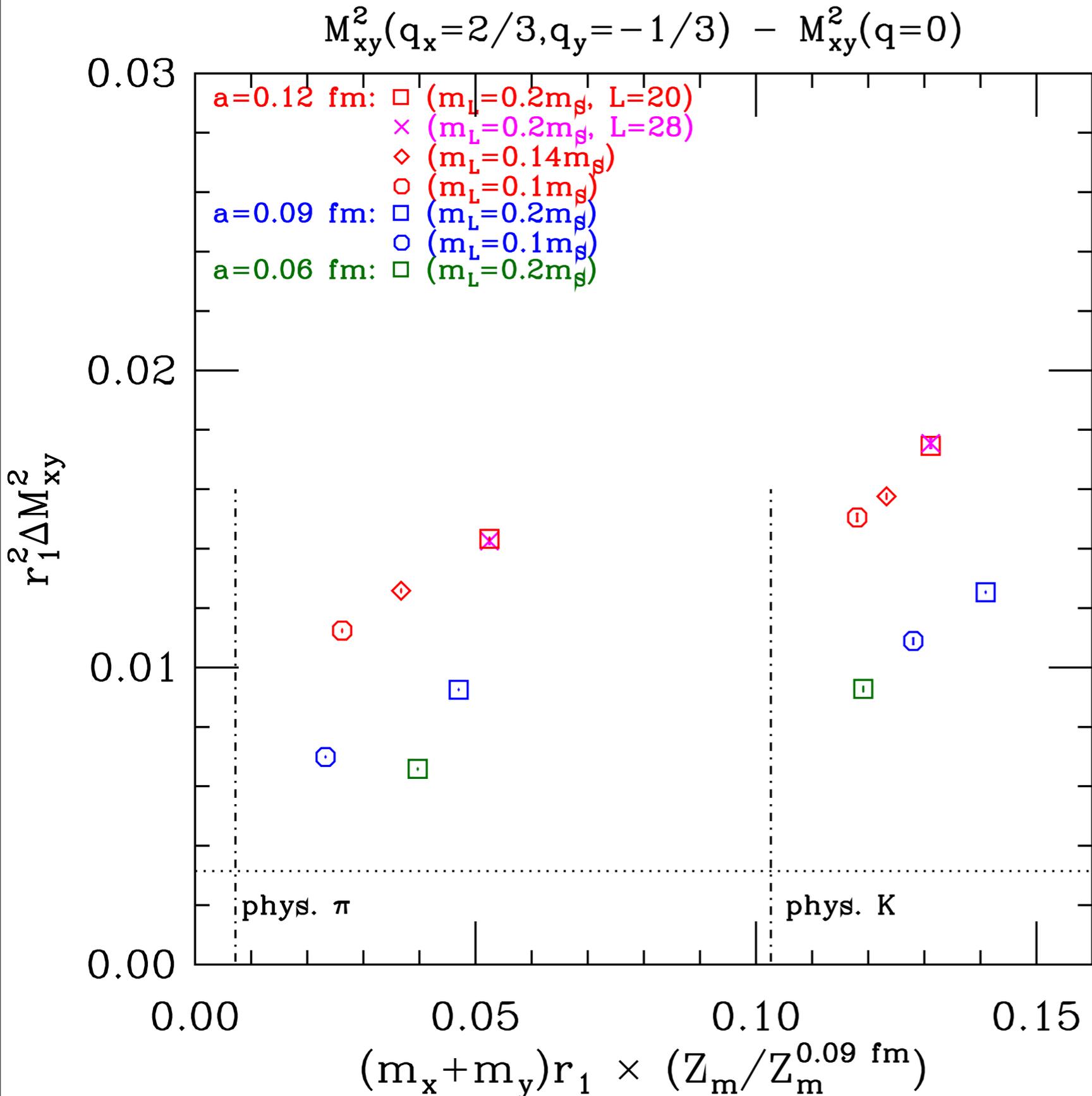
# FV Corrections: Comparison with Data

- ‘kaon’ and ‘pion’ points are the ones compared with BMW and RM123 results earlier.
- Each fit has 1 free parameter (overall height); shape is completely determined by ChPT at NLO.
- ChPT gives reasonable description of FV effects.
- Note that FV effect actually changes sign in ‘pion’ case.
- Can see why it is difficult to observe difference between results on  $L=20$  and  $L=28$  ensembles.



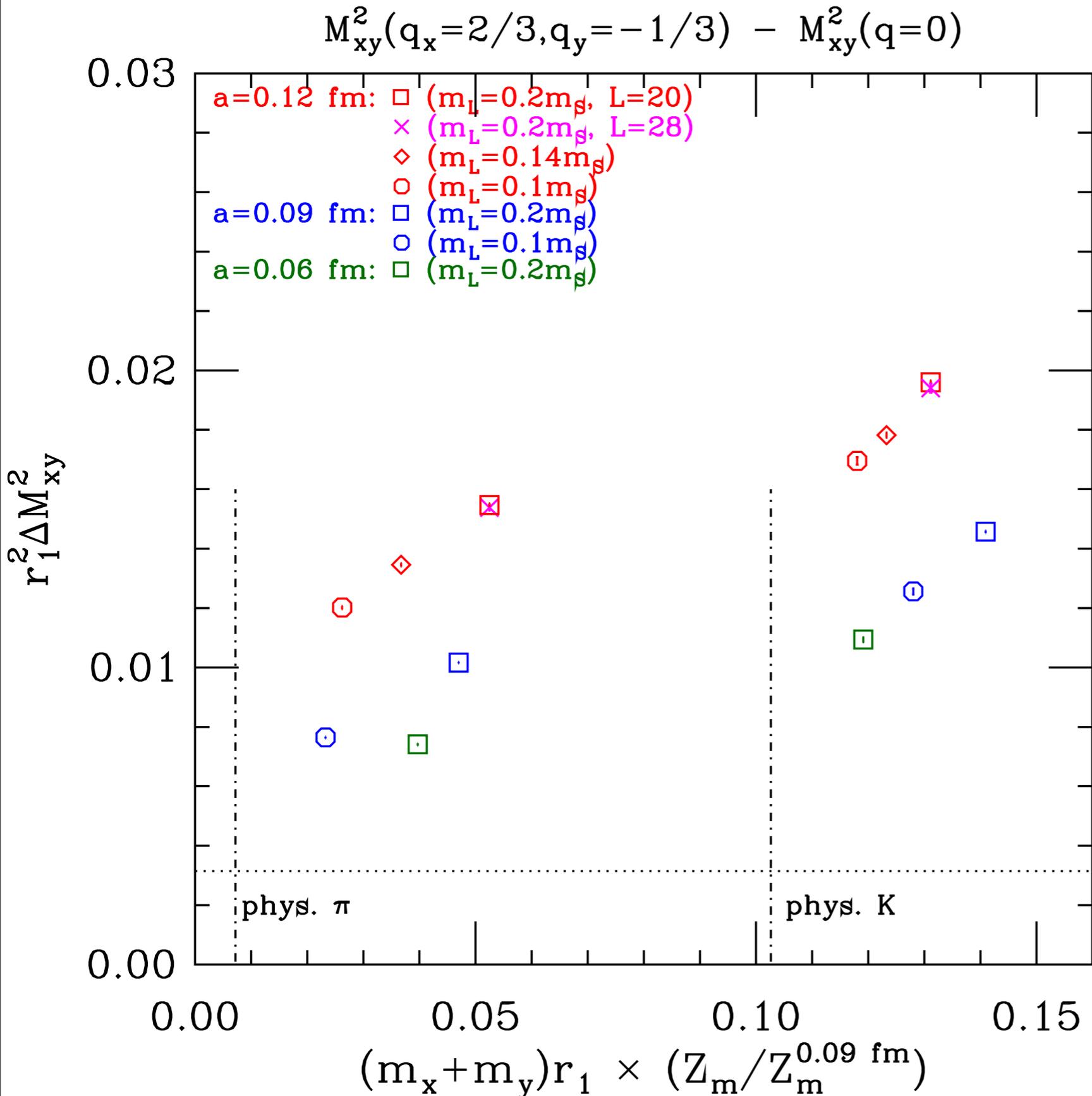
$a \approx 0.12$  fm,  $m_l/m_s = 0.01/.05$

# Chiral Fit and Extrapolation



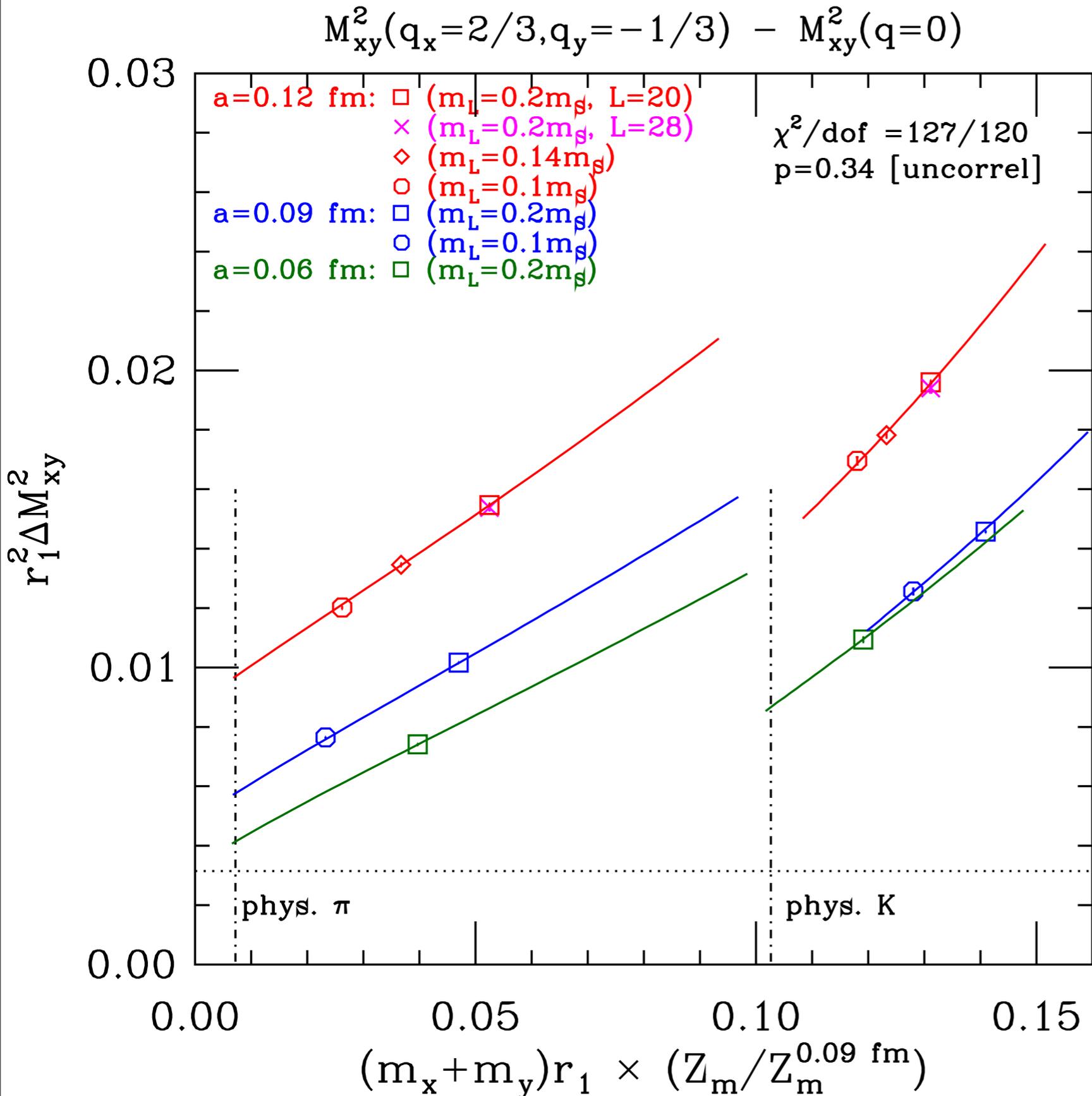
- Mass-square difference between charge +1 mesons ( $\pi^+$  &  $K^+$ ) and ones made from uncharged valence quarks
- Shows unitary points only.
- We have many partially quenched points, for charged and neutral mesons, as well as points with  $2 \times$  physical charges.
  - $\sim 150$  points in typical fit.
- A big part the difference between results from different lattice spacings is from mistuned  $m_S$ , not discretization effects.

# Chiral Fit and Extrapolation



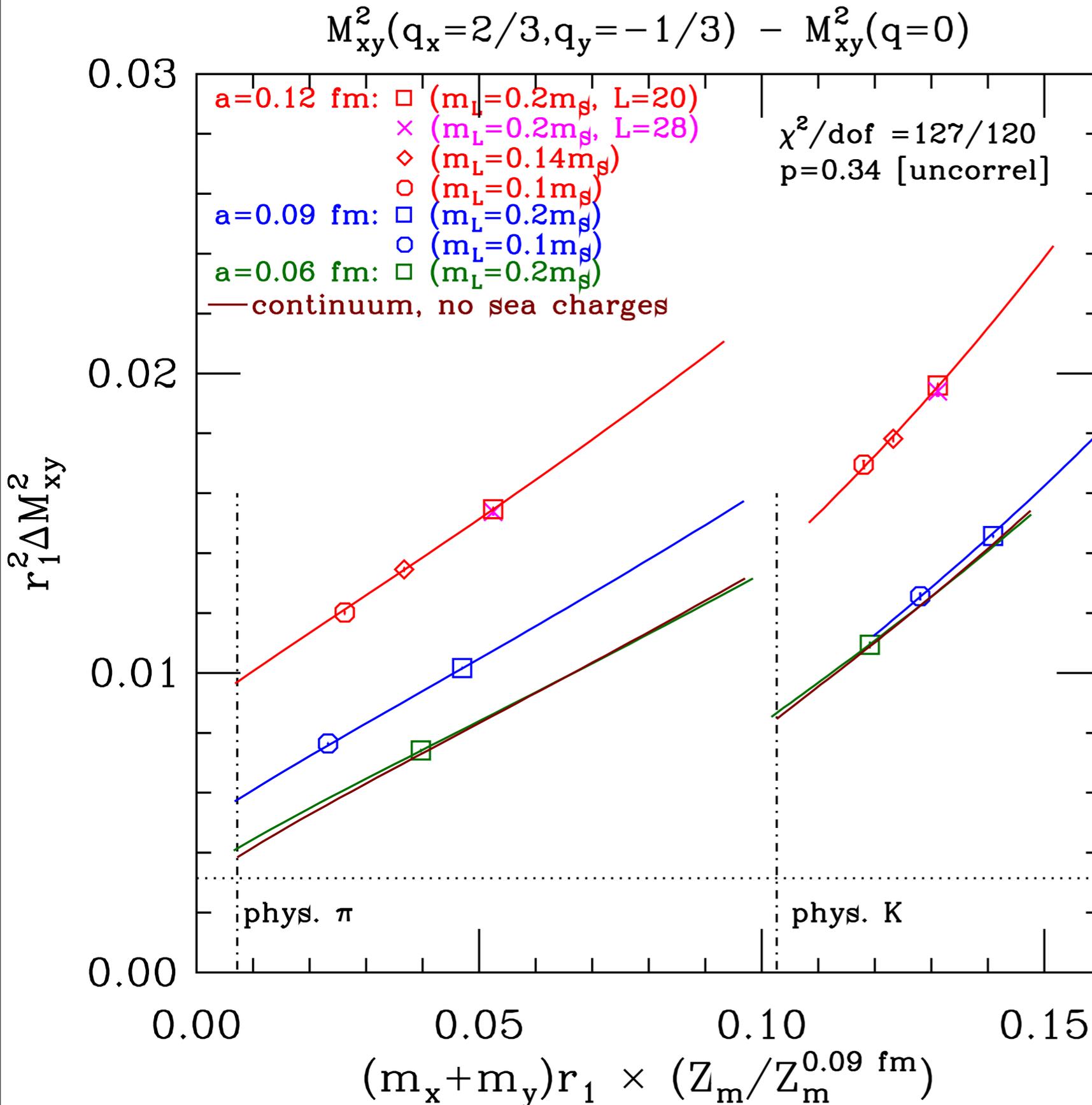
- Points after correction for finite-volume effects.
- Correction is  $\sim 7\text{--}10\%$  (pions) and  $\sim 10\text{--}18\%$  (kaons).
  - *Bigger* correction at higher mass because of overall factor of  $m^2$  in 1-loop diagrams, but not at LO (Dashen's theorem).
- Note that  $a \approx 0.12$  fm,  $m_l \approx 0.2m_s$  points for  $L=20$  ( $\square$ ) and  $L=28$  ( $\times$ ) are consistent.

# Chiral Fit and Extrapolation



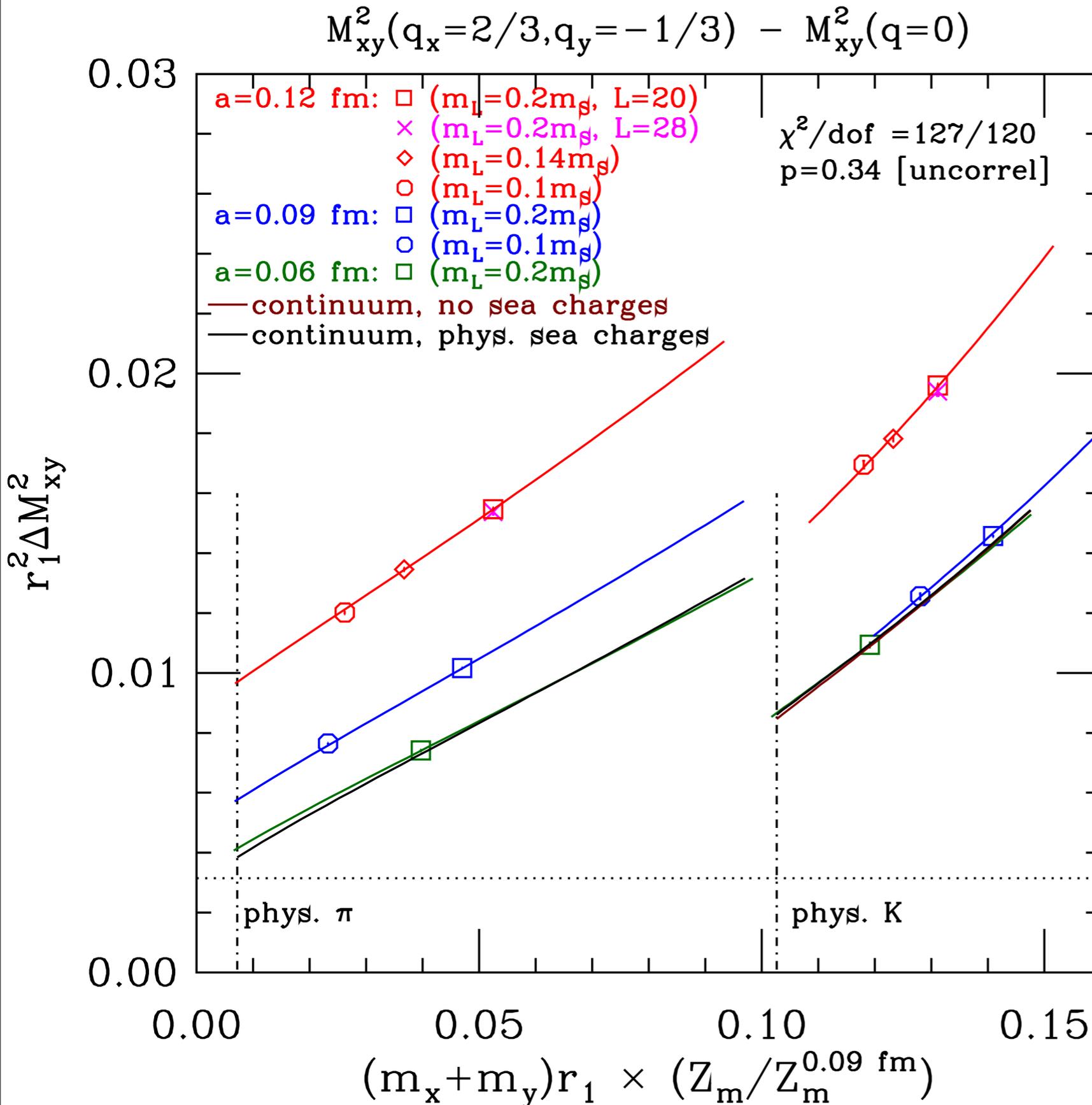
- Chiral fit to infinite-volume (corrected) points.
- Data has very high correlations for different valence masses or charges on the same ensembles: covariance matrix nearly singular.
- For that reason, and because errors are tiny (0.4--0.8%), it is difficult to get decent correlated fits.
- This is an uncorrelated fit; has 149 data points, 29 parameters,  $\chi^2/\text{dof}=127/120$ ,  $p=0.34$ .
- Fits are generally significantly better than earlier ones without FV corrections.

# Chiral Fit and Extrapolation



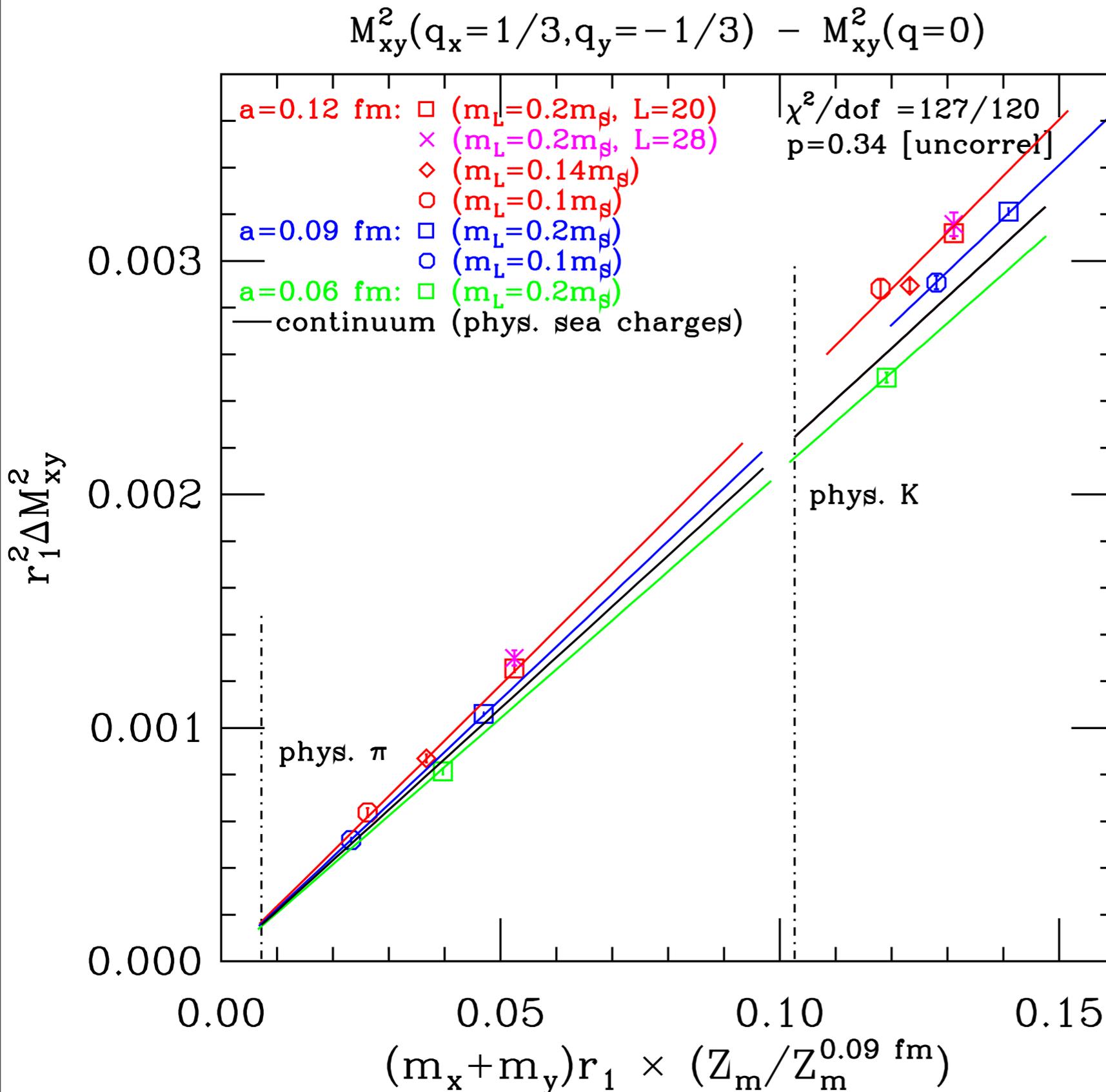
- Extrapolate to continuum, and set valence, sea masses equal.
- Adjust  $m_s$  to physical value.
- Keep sea charges = 0.
- Small change between  $a=0.06$  fm and continuum is conspiracy between discretization and  $m_s$  effects.

# Chiral Fit and Extrapolation



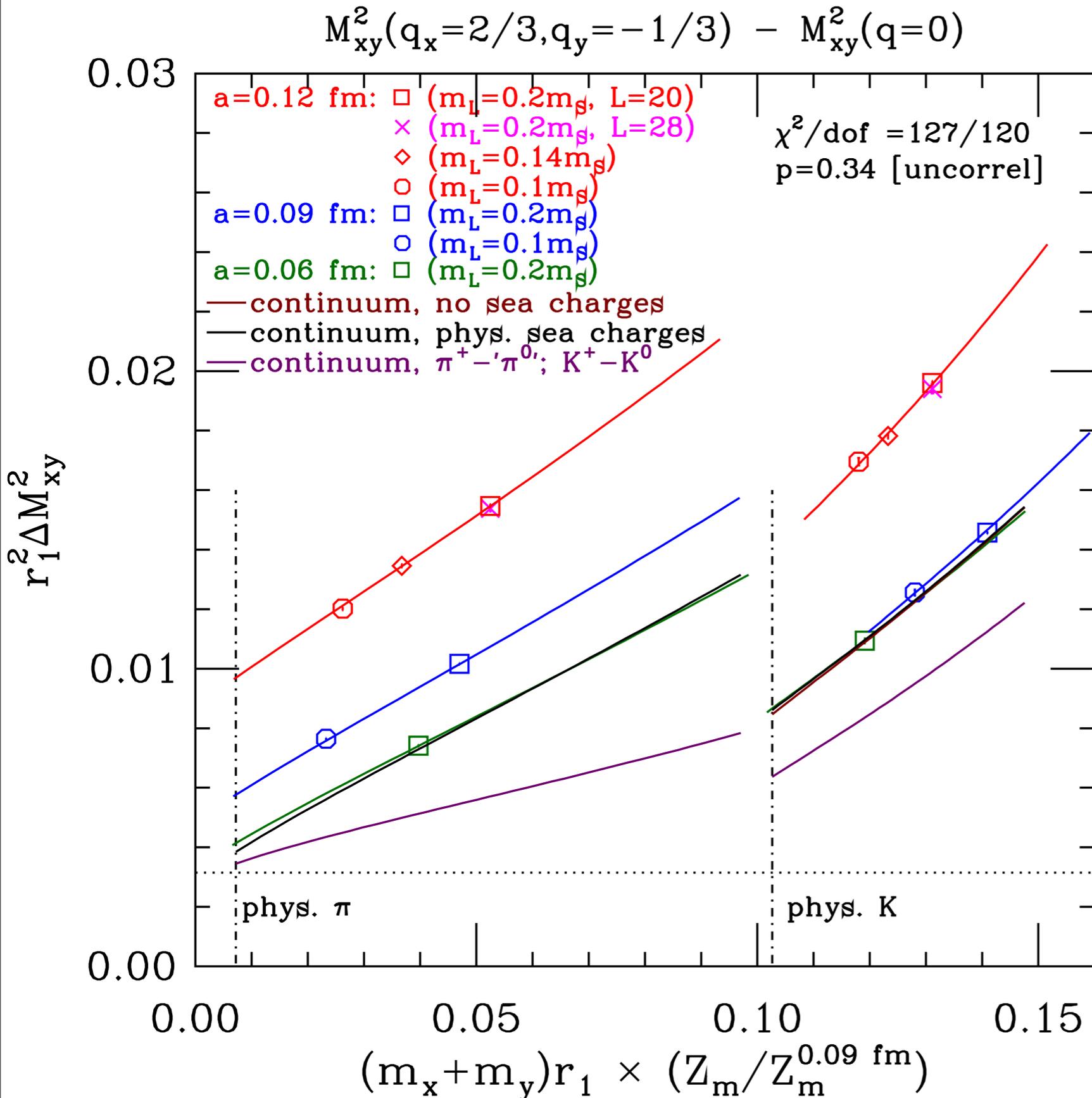
- Set sea quark charges to their physical values, using NLO chiral logs.
- Difference with previous case is very small for kaon; vanishes identically for pion.

# Chiral Fit and Extrapolation



- Neutral  $d\bar{d}$ -like mesons ( $q_x = q_y = 1/3$ ) for same fit.
- Note difference in scale from charged meson plot.
- $\sim$ Function of  $(m_x + m_y)$  only ( $\pi$  and  $K$  line up).
- Nearly linear: chiral logs vanish for neutrals.

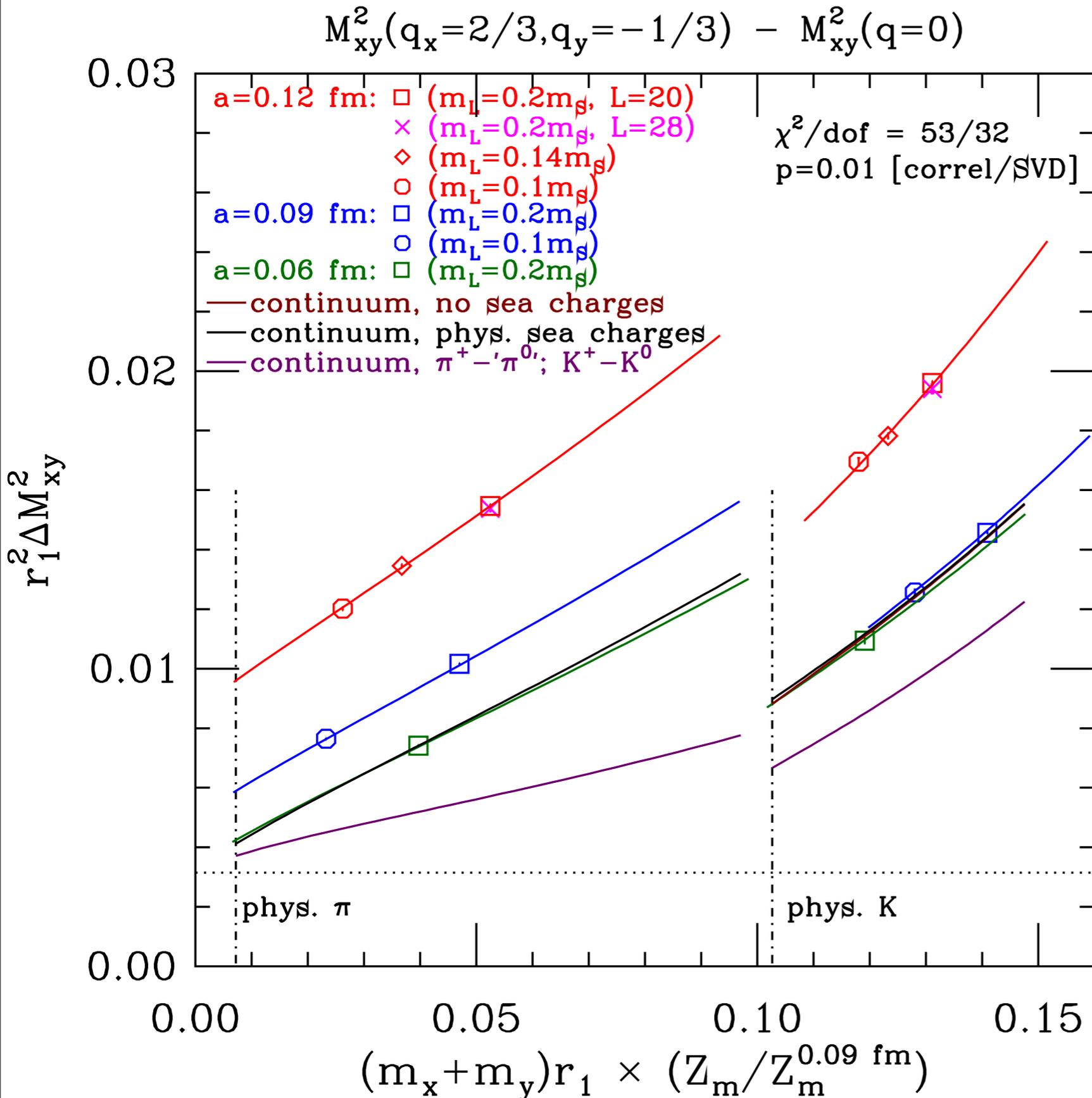
# Chiral Fit and Extrapolation



- Now subtract neutral masses from charged masses to give **purple** lines.
- We are not including disconnected EM graphs for  $\pi^0$ , which is why we call it ' $\pi^0$ '.
- Horizontal dotted line shows experimental value of  $\pi$  splitting; difference between it and intercept of **purple** line with vertical, dashed-dotted physical  $\pi$  line is a measure of systematic errors.
- Can now read off ratio of  $K$  and  $\pi$  splittings:

$$\epsilon = 0.84(5)$$

# Chiral Fit and Extrapolation



- Alternative correlated fit, with data that has been thinned more.
- SVD-like cut is needed; we cut eigenvalues of correlation matrix that are  $< 1$ .
- 55 data points, 23 params,  $\chi^2/\text{dof}=53/32$ ,  $p=0.01$ .
- Result is consistent with previous fit:

$$\epsilon = 0.79(8)$$

# Systematic Errors

- ◆ Difference between the finite-volume corrected result for  $\epsilon$  and the uncorrected one is 0.19. We currently take half this amount as the estimate of possible residual FV errors from higher orders in ChPT.
- ◆ Standard deviation on  $\epsilon$  over all current continuum/chiral fits is 0.13.
  - Here we include all uncorrelated fits with  $p > 10^{-3}$  or correlated fits with  $p > 10^{-8}$ .
- ◆ Instead of calculating  $\epsilon$  by ratio of results for  $K$  and  $\pi$  splittings, we may use the experimental  $\pi$  splitting. This gives  $\epsilon = 1.02(4)$ , or a difference from our central value of 0.18.
  - To be conservative, we take the larger number, 0.18, as an estimate of the lattice errors from the continuum/chiral extrapolation, although some of the difference may be due to residual finite-volume errors (included separately) or the effect of dropping disconnected diagrams for the  $\pi^0$ .

# Current Result

- ◆ Get (**preliminary!**):

$$\epsilon = 0.84(5)_{\text{stat}}(18)_{a^2}(10)_{\text{FV}}$$

or:

$$\epsilon = 0.84(21)$$

- ◆ Using this number with the current HISQ light meson analysis gives (**preliminary!**):

$$m_u/m_d = 0.4482(48)_{\text{stat}}\left(\begin{smallmatrix} + & 21 \\ - & 115 \end{smallmatrix}\right)_{a^2}(1)_{\text{FV}_{\text{QCD}}}(177)_{\text{EM}}$$

- where here “EM” denotes all errors from  $\epsilon$ , while “FV<sub>QCD</sub>” refers to finite-volume effects in the pure QCD calculation on the HISQ ensembles.

# Future Plans

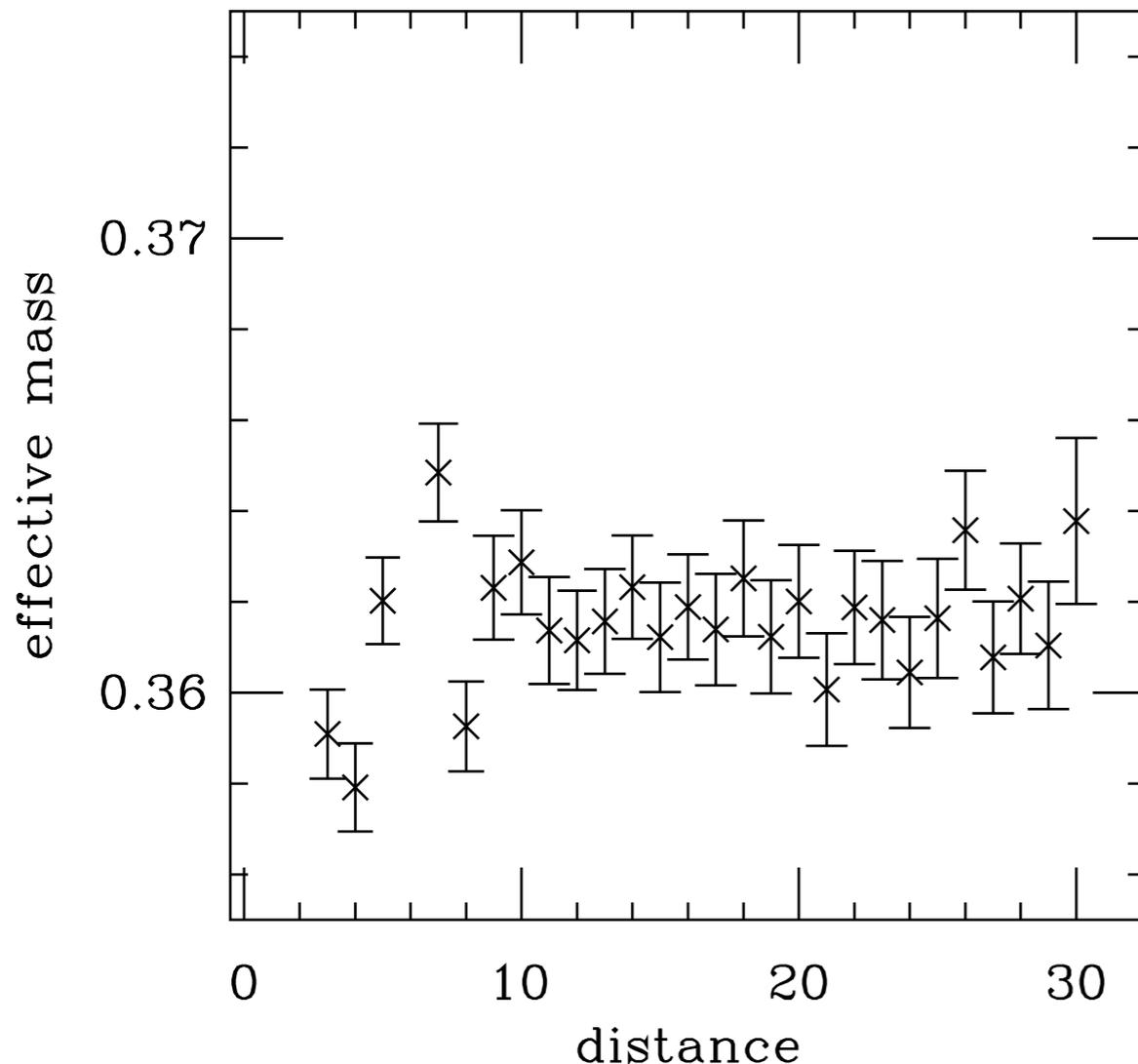
- ◆ We have data from additional ensembles at  $a \approx 0.06$  fm and  $a \approx 0.045$  fm.
  - need to complete analysis and add in to chiral/continuum extrapolation.
- ◆ EM effects in baryons also being studied.
- ◆ Extension to MILC HISQ ensembles is straightforward, and should reduce errors significantly:
  - Smaller discretization effects.
  - Nearly absent chiral extrapolation errors, since ensembles with physical masses are included.
  - Smaller FV effects, since our HISQ lattices are generally larger than the older asqtad ones. Max size  $\sim 5.5$  fm.
- ◆ Extension to unquenched case will make possible controlled calculations of many additional quantities.
  - Dynamic (unquenched) QED code has been written, and has passed some basic tests.

# Back-up Slides

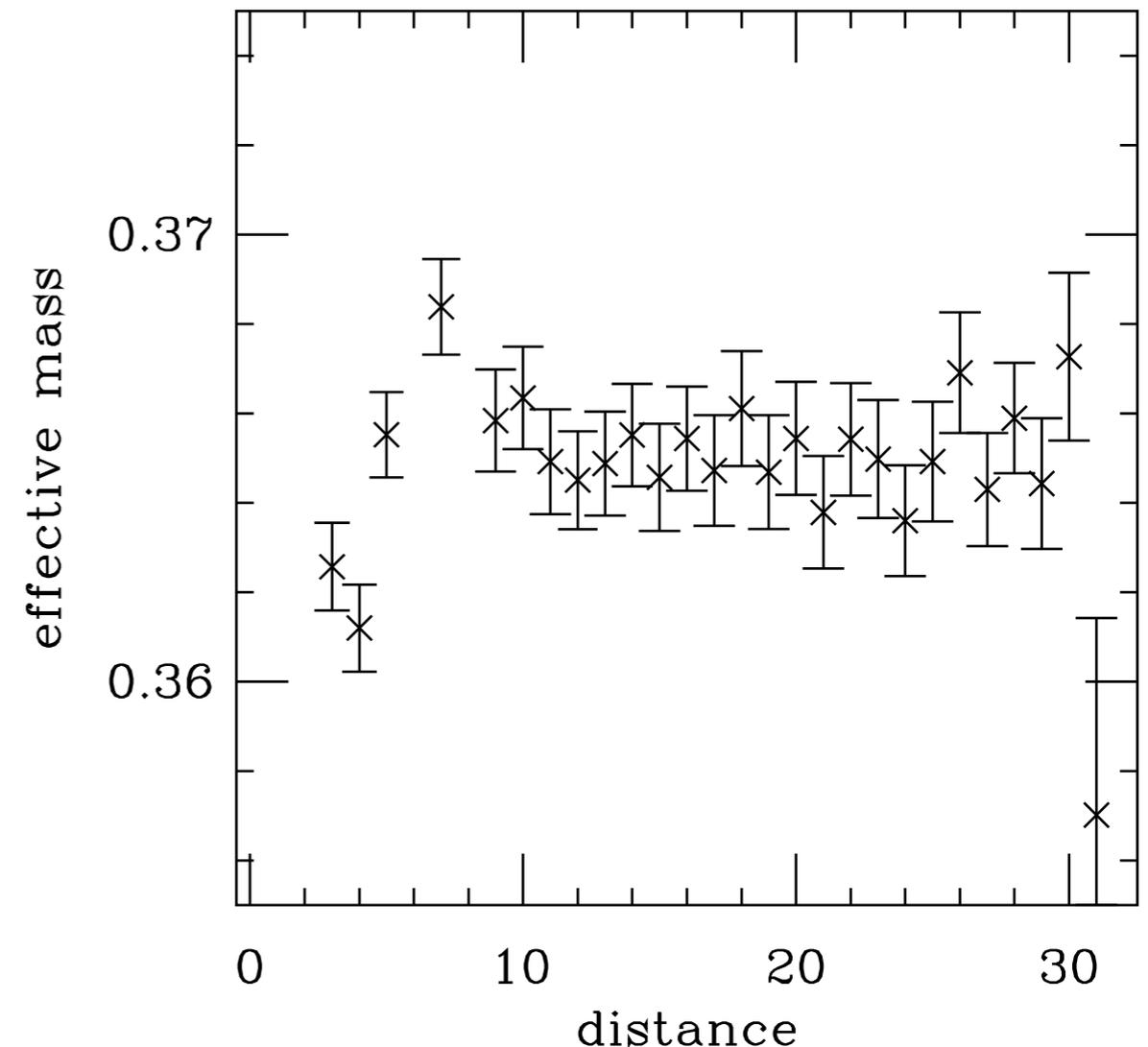
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# Effective Mass Plots

Kaon,  $12^3 \times 64$ , 0.01/.04,  $q=0$

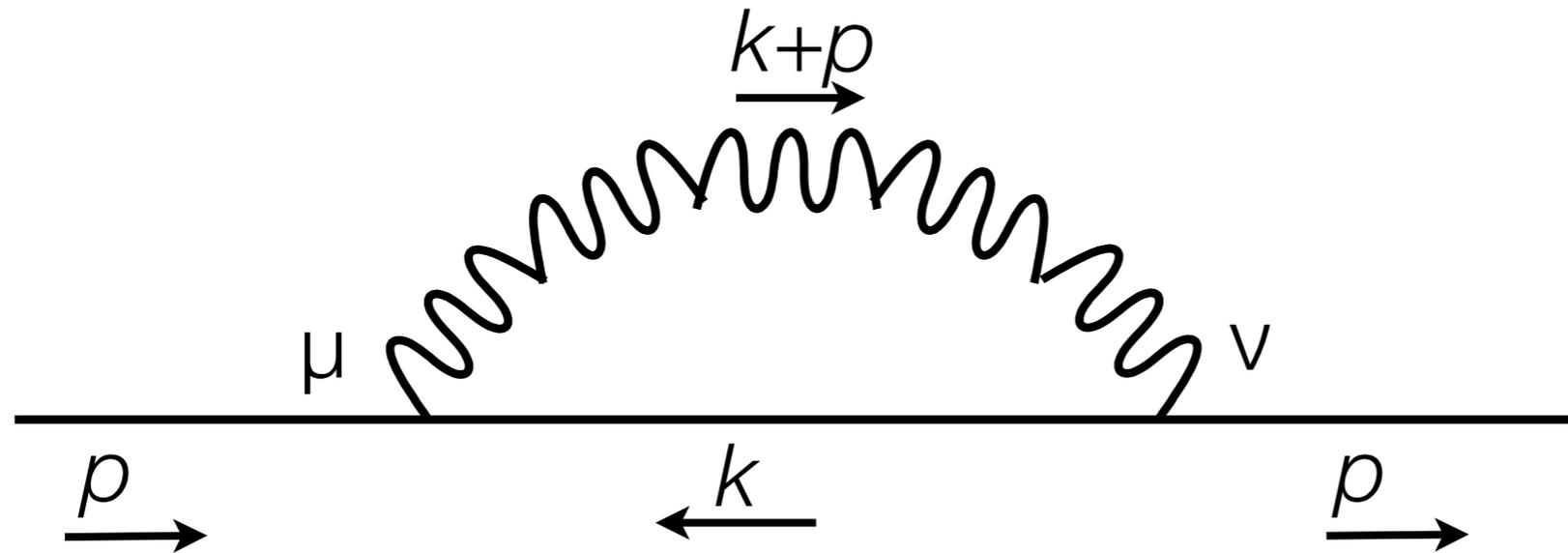


Kaon,  $12^3 \times 64$ , 0.01/.04,  $q=+1$



- No evidence of any systematic problem in extracting masses in charged case (right) compared to uncharged case (left).
- [BMW \[arXiv:1406.4088\]](https://arxiv.org/abs/1406.4088) recently reported problems (close excited states) in extracting masses for the FV version of EM that we use, but in pure (quenched) QED.
- We see no such problems in our quenched QED + full QCD simulations. But agree that masses are  $T$ -dependent, as seen in the FV formulas.

# Finite-Volume ChPT: Sunset Graph



- ◆ In rest frame,  $p = (p_0, 0, 0, 0)$ , only the 00 component of the photon propagator contributes.
- ◆ In infinite-volume, get:

$$\int \frac{d^4 k}{(2\pi)^4} \frac{1}{\vec{k}^2} \frac{(2p_0 + k_0)^2}{k^2 + m^2}$$

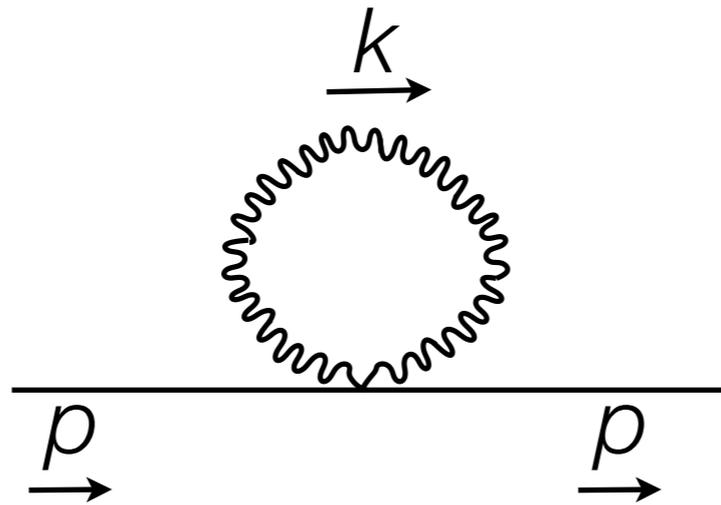
- where  $m$  is the meson mass, and numerator comes from momentum factors in the coupling of a (pseudo)scalar particle to a photon.

# Finite-Volume ChPT: Sunset Graph

$$\int \frac{d^4 k}{(2\pi)^4} \frac{1}{\vec{k}^2} \frac{(2p_0 + k_0)^2}{k^2 + m^2}$$

- ◆  $k_0$  integral, by itself, is linearly divergent.
- ◆ Even when we take difference between finite (spatial) volume version [“FV”] and infinite (spatial) volume version [“ $\infty V$ ”], the  $k_0$  integral makes the difference linearly divergent.
- ◆ (Usually, all divergences are the same in FV and  $\infty V$ , so difference diagram by diagram is finite.)
- ◆ Problem here is coming from lack of Lorentz covariance of the gauge.
- ◆ But photon tadpole has a piece that cancels the spurious  $k_0$  divergence.

# Finite-Volume ChPT: Photon Tadpole



- ◆ 00 piece of photon propagator gives:  $-\int \frac{d^4 k}{(2\pi)^4} \frac{1}{\vec{k}^2}$
- ◆ Combines with sunset to give:

$$\int \frac{d^4 k}{(2\pi)^4} \frac{1}{\vec{k}^2} \left( \frac{(2p_0 + k_0)^2}{k^2 + m^2} - 1 \right)$$

- ◆ finite-volume effects of this integral (FV -  $\infty$ V) are now finite & calculable.
- ◆ Do by brute force difference of FV sum from  $\infty$ V integral.
  - FV sum over  $2\pi n_i/L$  for spatial directions;  $2\pi n_0/T$  for time direction.

# FV Corrections: Comparison

- Accidental very small FV difference between  $20^3 \times 64$  (magenta) and  $28^3 \times 64$  (black) lattices at RM123 comparison point.

